Stochastic Excitation of solar-like acoustic modes, and its impact on mode asymmetry

Jordan Philidet, LESIA

K. Belkacem, M.-J. Goupil, R. Samadi, C. Barban, H.-G. Ludwig



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Meudon



d'île-de-France

Laboratoire d'Études Spatiales et d'Instrumentation en Astrophysique

Context

What is asterosismology?

Waves are **reflected** at the surface and **refracted** in the interior



Credits: Cunha et al. (2007)

The star acts as a resonant cavity, within which resonant modes develop, which can be either stable or unstable

- These modes are observed at the surface, by computing the Fourier transform of long time series.



Solar spectrum observed by GOLF

Context

Asterosismology: plenty of data

We detect solar-like modes in a **large amount** of stars, at various evolutionary stages and throughout the HR diagram (CoRoT, Kepler, TESS)

Data will be even more plentiful will PLATO



Credits: Chaplin & Miglio (2013)

Context

Why do we do asterosismology?

The resonant frequencies give information on the **internal structure** of the star:

$$\nu_{n,l} = \Delta \nu \left(n + \frac{l}{2} + \widetilde{\alpha} \right) + \epsilon_{n,l}$$

...

with
$$\Delta\nu = \left(2\int_0^R \frac{dr}{c(r)}\right)^{-1} \propto \ \overline{\rho}^{1/2}$$

For solar-like oscillators, the amplitudes and linewidths give information on the **turbulent convection**:

- Sub-surface convection is very turbulent ($Re \sim 10^{14}$) \rightarrow stochastic processes
- Mode amplitude is the result of a balance between stochastic excitation and damping
- Mode linewidth is related to the rate at which the mode is stochastically damped



What is it ?

The line profile of the modes in the Fourier spectrum are **asymmetric**



What is it ?

log₁₀ Power (arbitrary units)

The first contribution to mode asymmetry is the **localisation of the source** (Duvall et al. 1993) The two waves propagating from the source in opposite directions interfere more The line profile of the constructively here than here modes in the Fourier spectrum are asymmetric 2.5 2 1.5 -.2 0 .2 -.1 .1 $\delta \nu$ (mHz)

Credits: Duvall et al. (1993)

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The **correlated background** affects this interference pattern: it is the second contribution to mode asymmetry (Nigam & Kosovichev, 1998)



Why is it important ?

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 - The correlated background affects this interference pattern: it is the second contribution to mode asymmetry (Nigam & Kosovichev, 1998)
 - Prior studies had simplified models (point-like source), very parametrised, and they fit each mode individually (Nigam et al. 1998, Kumar & Basu 1999, Severino et al. 2001, Toutain et al. 2006). Goal: mainly to infer the position of the source



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We need a more realistic model for the source, with very few parameters, to compare the results with observations. Goal: **to make quantitative predictions applicable to stars other than the Sun**









Main results

- The predictions of the model are in agreement with observations → model validated
- The spatial extent of the source is crucial (sensitivity to how the properties of turbulence change above and below the photosphere)
 - The contribution of the **correlated background is negligible**



Credits: Philidet et al. (submitted to A&A)

Asymmetry reversal What is it ?

Mode asymmetry have opposite sign in velocity and in intensity



Why is it important?

Mode asymmetry have opposite sign in velocity and in intensity

There is no consensus on the explanation:

- correlated background higher than in velocity (Nigam & Kosovichev 1998)
- radiative effects (Georgobiani et al. 2001)
- non-adiabatic effects (Duvall et al. 1993)
- \rightarrow still an unsolved problem
- \rightarrow up to now, no predictive model



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Credits: Nigam et al. 1998

Asymmetries in other stars are measured in intensity → adapting our model to intensity spectrum is a key step to **apply the model to stars other than the Sun**



Credits: Benomar et al. 2018

Adaptation to intensity spectrum

$$\frac{\mathrm{d}^2 \Psi_\omega}{\mathrm{d}r^2} + \left(\frac{\omega^2 + j\omega\Gamma}{c^2} - V(r)\right)\Psi_\omega = S(r)$$

 $\frac{\delta L}{L} = \frac{L_R}{L} \left(\underbrace{\left[(4+3\kappa_\rho)A_{\xi} + (4-\kappa_T)A_T + r\kappa_\rho \frac{\mathrm{d}A_{\xi}}{\mathrm{d}r} - H_T \frac{\mathrm{d}A_T}{\mathrm{d}r} \right]}_{\equiv A_L} \Psi_{\omega} \right)$ $+ \underbrace{\left[(4-\kappa_T)B_T + r\kappa_\rho A_{\xi} - H_T \left(A_T + \frac{\mathrm{d}B_T}{\mathrm{d}r}\right) \right]}_{\equiv B_L} \underbrace{\frac{\mathrm{d}\Psi_{\omega}}{\mathrm{d}r}}_{\equiv C_L} - H_T B_T \frac{\mathrm{d}^2\Psi_{\omega}}{\mathrm{d}r^2}}_{\equiv C_L} \right)$

The intensity eigenfunction do not depend only on Ψ_{ω} , but also on its radial derivatives The nodes in velocity and intensity are different

The asymmetries are different too

Main results

- We find agreement with observations with the same model and same parameter values as in velocity → no additional mechanism needed
- Radiative effects reduce the amplitude of the mode in intensity, but do not affect asymmetry
- Discrepancies at high frequency in intensity:
 Non-adiabaticity ?
 - Correlated background ?
 - Contribution of convective flux ?



Credits: Philidet et al. (in prep)

Thank you for your attention !